

## Lecture 17

Thursday, November 14, 2024 11:36

# Chapter 10 : Electron Spin and the Spin-Statistics Theorem

## 10.1 Electron Spin

Picture the electron as a sphere of charge spinning about one of its diameters.



note: electron spin has no classical analogue

Electrons have intrinsic angular momentum.

Spin angular momentum or spin.

$$\hat{S}^2 = \hat{S}_x^2 + \hat{S}_y^2 + \hat{S}_z^2$$

## Spin angular momentum operators

$$[\hat{L}_x, \hat{L}_y] = i\hbar \hat{L}_z \rightarrow [\hat{S}_x, \hat{S}_y] = i\hbar \hat{S}_z$$

$$[\hat{L}_y, \hat{L}_z] = i\hbar \hat{L}_x \rightarrow [\hat{S}_y, \hat{S}_z] = i\hbar \hat{S}_x$$

$$[\hat{L}_z, \hat{L}_x] = i\hbar \hat{L}_y \rightarrow [\hat{S}_z, \hat{S}_x] = i\hbar \hat{S}_y$$

$$[\hat{S}^2, \hat{S}_x] = 0, [\hat{S}^2, \hat{S}_y] = 0, [\hat{S}^2, \hat{S}_z] = 0$$

Eigenvalues of  $\hat{S}^2$  :

$$s(s+1)\hbar^2, s = 0, \frac{1}{2}, 1, \frac{3}{2}, \dots$$

Eigenvalues of  $\hat{S}_z$  :

$$m_s \hbar, m_s = -s, -s+1, \dots, s-1, s$$

Quantum number  $s$  : spin.

Electrons can only have  $s = \frac{1}{2}$

Protons, neutrons :  $s = \frac{1}{2}$

Photons :  $s = 1$

(exception : photons can have  $m_s = -1$  or  $m_s = 1$ , but not  $m_s = 0$ )

$$s = \frac{1}{2}$$

$$\left[ \left(\frac{1}{2}\right) \left(\frac{3}{2}\right) \hbar^2 \right]^{\frac{1}{2}} = \frac{1}{2} \sqrt{3} \hbar$$

for  $\hat{S}_z$ ,  $m_s = +\frac{1}{2}\hbar$  and  $m_s = -\frac{1}{2}\hbar$

$$\hat{S}_z \alpha = +\frac{1}{2} \hbar \alpha$$

$$\hat{S}_z \beta = -\frac{1}{2} \hbar \beta$$

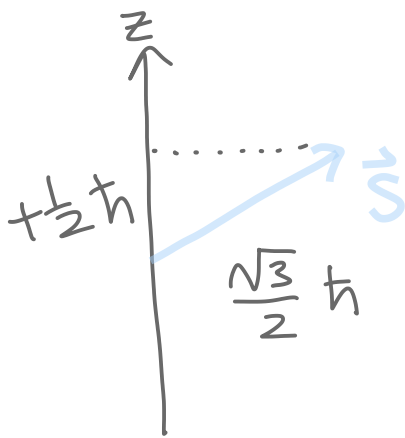
$\alpha, \beta$  are  
eigenfunctions  
of  $\hat{S}_z$

Because  $\hat{S}_z$  commutes with  $\hat{S}^2$ ,

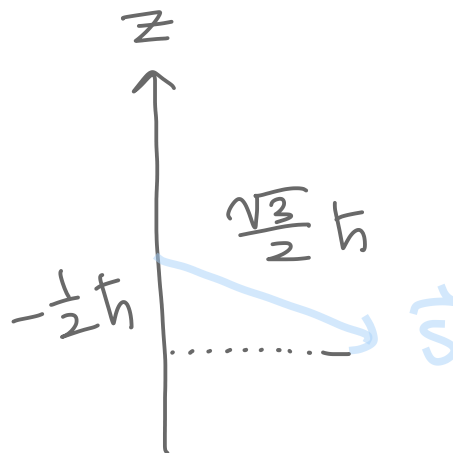
$\alpha$  and  $\beta$  are also eigen functions of  $\hat{S}^2$ .

$$\hat{S}^2 \alpha = \frac{3}{4} \hbar^2 \alpha$$

$$\hat{S}^2 \beta = \frac{3}{4} \hbar^2 \beta$$



spin up



spin down

$$\Psi(x, y, z) \alpha(m_s)$$

$$\Psi(x, y, z) \beta(m_s)$$

$$\sum_{m_s = -\frac{1}{2}}^{\frac{1}{2}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} |\Psi(x, y, z, m_s)|^2 dx dy dz = 1$$

## 10.2 Spin and the Hydrogen Atom

$$\Psi(x, y, z) g(m_s)$$

where  $g(m_s)$  is either  $\alpha$  or  $\beta$ .

more generally,

$$g(m_s) = C_1 \alpha + C_2 \beta$$

$$\hat{H} [\Psi(x, y, z) g(m_s)]$$

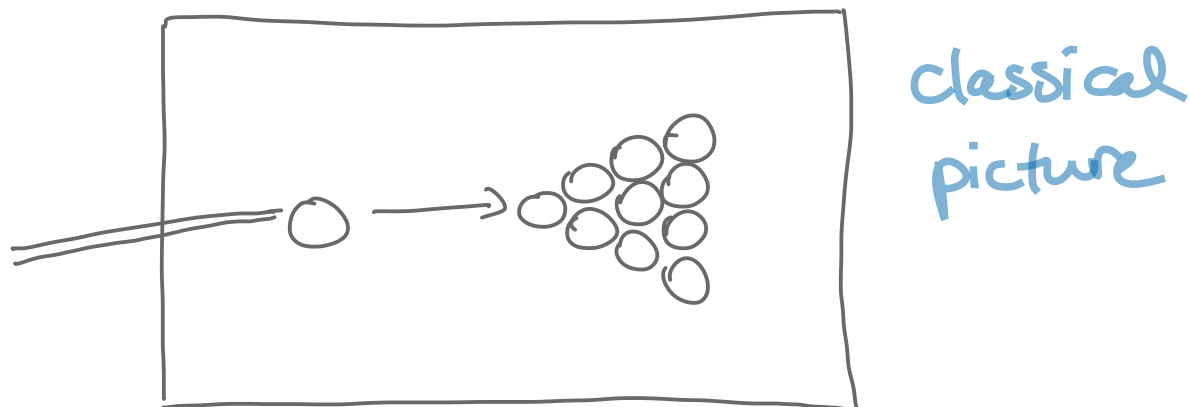
$$= g(m_s) \hat{H} \Psi(x, y, z)$$

$$= E [\Psi(x, y, z) g(m_s)]$$

$$\Psi(x, y, z) \alpha \quad \text{or} \quad \Psi(x, y, z) \beta$$

Degeneracy of H-atom energy levels becomes  $2n^2$  rather than  $n^2$ .

## 10.3 The Spin-Statistics Theorem



Even if billiard balls were identical, we can still trace their individual paths on a macroscopic level.

For microscopic particles, it is not possible to trace their paths because of the uncertainty principle.

indistinguishability of  
electrons in the wavefunction

Consider a system of  $n$  particles:

particle  $l$  :  $x_l, y_l, z_l, m_{s_l}$

... to all four variables

Let  $q_i$  denote all the variables  
 ( $q$  is typically used to denote  
 generalized coordinates)

Exchange or permutation operator  $\hat{P}_{12}$   
 interchanges all the coordinates of  
 particles 1 and 2.

$$\hat{P}_{12} \Psi(q_1, q_2, q_3, \dots, q_n) = \Psi(q_2, q_1, q_3, \dots, q_n)$$

Example :

$$\hat{P}_{12} [s(1)\alpha(1)\beta(2)] = s(2)\alpha(2)\beta(1)$$

What are the eigenvalues of  $\hat{P}_{12}$ ?

$$\begin{aligned} & \hat{P}_{12} \hat{P}_{12} f(q_1, q_2, \dots, q_n) \\ &= \hat{P}_{12} f(q_2, q_1, \dots, q_n) \end{aligned}$$

$$= f(q_1, q_2, \dots, q_n)$$

applying  $\hat{P}_{12}$  twice has no net effect

$$\hat{P}_{12}^2 = \hat{I}$$

Let  $c_i$  and  $w_i$  denote the eigenvalues and eigenfunctions of  $\hat{P}_{12}$ .

$$\hat{P}_{12} w_i = c_i w_i$$

$$\hat{P}_{12}^2 w_i = c_i \hat{P}_{12} w_i$$

$$\hat{I} w_i = c_i^2 w_i$$

$$c_i = \pm 1$$

$$\hat{P}_{12} w_+(q_1, q_2, \dots, q_n) = (+1) w_+(q_1, q_2, \dots, q_n)$$

symmetric



symmetric

$$\hat{P}_{12} \omega_-(q_1, q_2, \dots, q_n) = (-1) \omega_-(q_1, q_2, \dots, q_n)$$

antisymmetric

$$\hat{P}_{ik} f(q_1, \dots, q_i, \dots, q_k, \dots, q_n)$$

$$= f(q_1, \dots, q_k, \dots, q_i, \dots, q_n)$$

$$\hat{P}_{ik} \Psi(q_1, \dots, q_i, \dots, q_k, \dots, q_n)$$

$$= c \Psi(q_1, \dots, q_i, \dots, q_k, \dots, q_n)$$

because the particles are indistinguishable

$\Psi$  is an eigenfunction of  $\hat{P}_{ik}$

$\therefore$  The wavefunction for a system

of  $n$  particles must be either symmetric or antisymmetric with respect to interchange of any two of the identical particles.

For electrons, the antisymmetric rule applies.

More generally, particles with half-integer spins ( $S = \frac{1}{2}, \frac{3}{2}, \dots$ ) require antisymmetric wavefunction  
fermions

Particles with integer spins ( $S = 0, 1, 2, \dots$ ) require symmetric wavefunctions.  
bosons

# Spin-Statistics Theorem

$$\Psi(q_1, q_2, \dots, q_n) = -\Psi(q_2, q_1, \dots, q_n)$$

What if electrons 1 and 2 have the same coordinates?

$$\text{i.e. } x_1 = x_2$$

$$y_1 = y_2$$

$$z_1 = z_2$$

$$m_{s_1} = m_{s_2}$$

$$q_1 = q_2$$

$$\Psi(q_1, q_1, \dots, q_n) = -\Psi(q_1, q_1, \dots, q_n)$$

$$2\Psi = 0$$

$$\Psi(q_1, q_1, \dots, q_n) = 0$$

Two electrons with the same  $s_f$  have zero probability of being found at the same point in space.



Pauli exclusion

For bosons, no such restriction

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## 10.4 The Helium Atom

For 2 electrons, we have

$$\alpha(1) \alpha(2)$$

$$\beta(1) \beta(2)$$

$$\alpha(1) \beta(2) \quad \left. \begin{array}{l} \alpha(2) \beta(1) \end{array} \right\} \text{violate indistinguishability}$$

$$\alpha(2)\beta(1)$$

Symmetric

$$\begin{aligned} & \alpha(1)\alpha(2) \\ & \beta(1)\beta(2) \\ & [\alpha(1)\beta(2) + \beta(1)\alpha(2)] 2^{-1} \end{aligned}$$

antisymmetric

$$[\alpha(1)\beta(2) - \beta(1)\alpha(2)] 2^{-1}$$